

Physics 517
 Fall 2004
 Due: 6 October 2004

Assignment # 5

1. The interval for Minkowski space can be written

$$ds^2 = -dt^2 + dx^2 + dy^2 + dz^2. \quad (1)$$

(We choose units where c , the speed of light, is unity, $c = 1$.) The Lorentz transformations are important in relativity and E&M because they leave the form of the interval invariant. Consider a frame \mathcal{O} with coordinates $\{t, x, y, z\}$, and a frame $\bar{\mathcal{O}}$ with coordinates $\{\bar{t}, \bar{x}, \bar{y}, \bar{z}\}$, moving with velocity v along the x -axis with respect to frame \mathcal{O} . The Lorentz transformation between these two frames can be written

$$\bar{t} = \gamma(t - vx), \quad (2)$$

$$\bar{x} = \gamma(x - vt), \quad (3)$$

$$\bar{y} = y, \quad (4)$$

$$\bar{z} = z, \quad (5)$$

where v is the boost velocity along the x -axis and $\gamma = (1 - v^2)^{-1/2}$.

- (a) Find a matrix representation of the transformation $\Lambda^{\bar{a}}_b$ and its inverse, $\Lambda^a_{\bar{b}}$.
 - (b) Write the metric g_{ab} in matrix form.
 - (c) Transform the metric to obtain $g_{\bar{a}\bar{b}}$, and write the solution in matrix form.
 - (d) The Riemann curvature tensor is often written R^a_{bcd} . Write an expression for $R^{\bar{a}}_{\bar{b}\bar{c}\bar{d}}$ in terms of R^a_{bcd} and the transformations matrices $\Lambda^{\bar{a}}_b$ and $\Lambda^a_{\bar{b}}$. Do not evaluate this expression.
2. \mathbf{v} and \mathbf{w} are vectors in the vector space, \mathcal{V} , with an orthonormal basis $\{\mathbf{e}_i\}_{i=1}^N$. The dual basis in \mathcal{V}^* is $\{\boldsymbol{\omega}^i\}_{i=1}^N$, and $\boldsymbol{\omega}^i(\mathbf{e}_j) = \delta^i_j$. A symmetric tensor \mathbf{T} is constructed from \mathbf{v} and \mathbf{w} as $\mathbf{T} = \frac{1}{2}(\mathbf{v} \otimes \mathbf{w} + \mathbf{w} \otimes \mathbf{v})(\boldsymbol{\theta}^1, \boldsymbol{\theta}^2)$, where $\boldsymbol{\theta}^1$ and $\boldsymbol{\theta}^2$ are elements of \mathcal{V}^* .
- (a) Expand \mathbf{v} and \mathbf{w} in the basis \mathbf{e}_i , and find the components of \mathbf{T} in this basis, T^{ab} , in terms of v^i and w^j . The components are found by letting \mathbf{T} act on the basis forms, $\boldsymbol{\omega}^i$: $T^{ab} = \mathbf{T}(\boldsymbol{\omega}^a, \boldsymbol{\omega}^b)$. Show that the contravariant components, T^{ab} , are symmetric, $T^{ab} = T^{ba}$.
 - (b) Show that the covariant components, T_{ab} , are symmetric, $T_{ab} = T_{ba}$. Begin with the result of part (a), and assume the existence of a metric, g_{ab} , that may be used to lower indices.
 - (c) \mathbf{A} is an anti-symmetric tensor, $A_{ab} = -A_{ba}$. Show that $A_{ab}T^{ab} = 0$.

3. \mathbf{A} is a generic rank (0,2) tensor with components A_{ab} .
- Write an expression for the completely anti-symmetric tensor $A_{[ab]}$ in terms of the components A_{ab} .
 - Write an expression for the completely symmetric tensor $A_{(ab)}$ in terms of the components A_{ab} .
 - Show that \mathbf{A} can be written in terms of $A_{(ab)}$ and $A_{[ab]}$.
 - \mathbf{C} is a symmetric tensor, $C^{ij} = C^{(ij)}$. Show that $A_{ij}C^{ij} = A_{(ij)}C^{ij}$.
 - \mathbf{D} is an anti-symmetric tensor, $D^{ij} = D^{[ij]}$. Show that $A_{ij}D^{ij} = A_{[ij]}D^{ij}$.
4. $\mathbf{v}_1, \mathbf{v}_2, \mathbf{v}_3$ and \mathbf{v}_4 , are elements of a 4-dimensional vector space, \mathcal{V} , with an orthonormal basis $\{\mathbf{e}_i\}_{i=1}^4$:

$$\mathbf{v}_1 = \mathbf{e}_1 + 3\mathbf{e}_2 - 2\mathbf{e}_3 + \mathbf{e}_4 \quad (6)$$

$$\mathbf{v}_2 = -\mathbf{e}_2 + 2\mathbf{e}_4 \quad (7)$$

$$\mathbf{v}_3 = \mathbf{e}_1 + 5\mathbf{e}_3 - 3\mathbf{e}_4 \quad (8)$$

$$\mathbf{v}_4 = -\mathbf{e}_1 - 12\mathbf{e}_3 + 13\mathbf{e}_4 \quad (9)$$

Are these vectors linearly independent? Answer this question by calculating the exterior product $\mathbf{v}_1 \wedge \mathbf{v}_2 \wedge \mathbf{v}_3 \wedge \mathbf{v}_4$. If the vectors are linearly independent, then they could form a basis in \mathcal{V} , and the exterior product must be non-zero. (Hint: Consider first calculating $\mathbf{v}_1 \wedge \mathbf{v}_2$ and $\mathbf{v}_3 \wedge \mathbf{v}_4$ individually, before calculating the final product. Why does this simplify the calculation?)

5. Show that when there is a sum over an upper index and a lower index, swapping the upper index to a lower index, and vice versa, does not change the sum. In other words, $A^a B_a = A_a B^a$.
6. $\{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}$ is an oriented, orthonormal basis in three dimensions. Evaluate these expressions with the Hodge star operator:
- $*1$
 - $*\mathbf{e}_1$
 - $*(\mathbf{e}_1 \wedge \mathbf{e}_3)$
 - $*(\mathbf{e}_1 \wedge \mathbf{e}_2 \wedge \mathbf{e}_3)$

7. The line element for a black hole spacetime can be written

$$ds^2 = -A^2 dt^2 + 2B dt dr + C^2 dr^2 + r^2(d\theta^2 + \sin^2 \theta d\phi^2), \quad (10)$$

where A , B and C are functions, and the coordinates are $\{t, r, \theta, \phi\}$.

- (a) Find the metric components g_{ab} . Write the metric in matrix form.
 - (b) The four-velocity of a particle in this spacetime is $u^a = (1, 0, 0, 0)$. Find the components of u_a .
8. \mathbf{u} , \mathbf{v} and \mathbf{w} are vectors in 3-dimensional Euclidian space. Show that
- (a) $\mathbf{*}(\mathbf{u} \wedge \mathbf{v})$ is equivalent to $\mathbf{u} \times \mathbf{v}$;
 - (b) $\mathbf{*}(\mathbf{*}\mathbf{u} \wedge \mathbf{v})$ is equivalent to $\mathbf{u} \cdot \mathbf{v}$;
 - (c) $\mathbf{*}(\mathbf{w} \wedge \mathbf{u} \wedge \mathbf{v})$ is equivalent to $\mathbf{w} \cdot (\mathbf{u} \times \mathbf{v})$;
 - (d) $\mathbf{*}(\mathbf{w} \wedge \mathbf{*}(\mathbf{u} \wedge \mathbf{v}))$ is equivalent to $\mathbf{w} \times (\mathbf{u} \times \mathbf{v})$.
9. Let \mathbf{u} be an unit vector, $\mathbf{u} \cdot \mathbf{u} = 1$. A projection operator, \mathbf{P}_u , may be defined by

$$\mathbf{P}_u = \mathbf{g} - \mathbf{u} \otimes \mathbf{u} \quad (11)$$

where \mathbf{g} is the metric tensor. \mathbf{P}_u projects an arbitrary vector \mathbf{v} into \mathbf{v}_\perp , which is orthogonal to \mathbf{u} : $\mathbf{v}_\perp = \mathbf{P}_u(\mathbf{v})$ and $\mathbf{v}_\perp \cdot \mathbf{u} = 0$. In component form, the projection can be written

$$v_\perp^a = P^a_b v^b = (g^a_b - u^a u_b) v^b \quad (12)$$

- (a) Show explicitly that $\mathbf{v}_\perp \cdot \mathbf{u} = 0$.
 - (b) Show that multiple applications of \mathbf{P}_u do not change the vector \mathbf{v}_\perp , i.e., $\mathbf{P}_u(\mathbf{v}_\perp) = \mathbf{v}_\perp$.
 - (c) Show that \mathbf{P}_u is the metric tensor for vectors perpendicular to \mathbf{u} : $\mathbf{P}_u(\mathbf{v}_\perp, \mathbf{w}_\perp) = \mathbf{g}(\mathbf{v}_\perp, \mathbf{w}_\perp) = \mathbf{v}_\perp \cdot \mathbf{w}_\perp$. (Write \mathbf{P}_u with both indices down.)
10. **Either** work problem 2.5.10 from Arfken, **or** explain in two sentences (or fewer) why you would rather not work problem 2.5.10 from Arfken.